

Lecture 9

Green's Functions for ODEs and PDEs

References

Haberman APDE, Sec. 8.2 and 8.3.

Haberman APDE, Sec. 9.1 and 9.2.

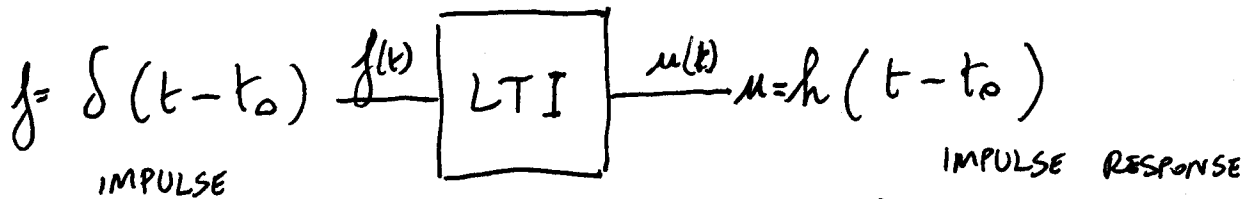
http://en.wikipedia.org/wiki/Green's_function

See also Lecture 8 notes.

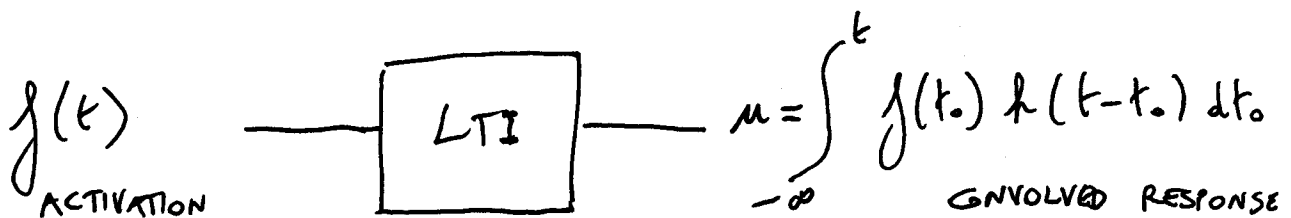
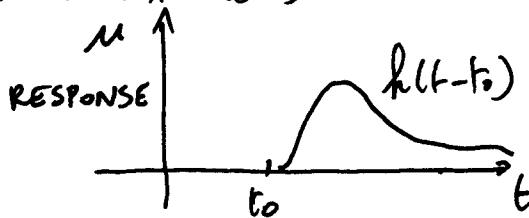
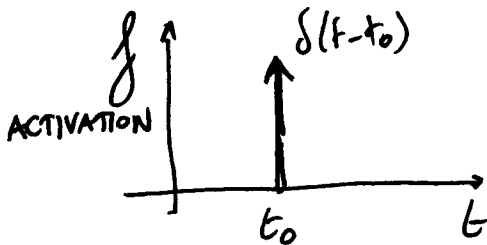
GREEN'S FUNCTION FOR ODE S

$$\mathcal{L} u = f$$

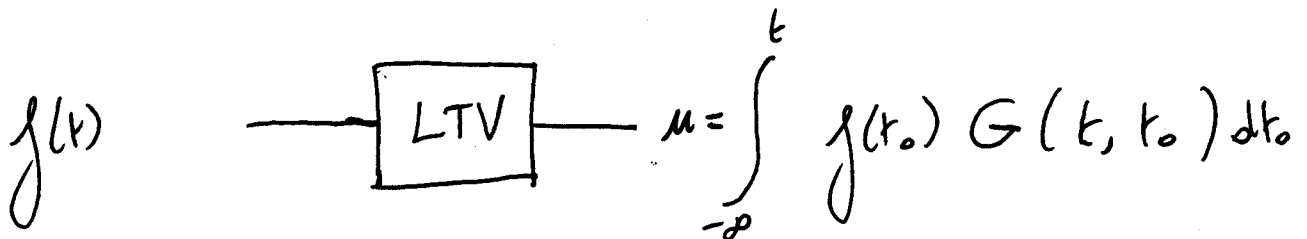
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 RESPONSE ACTIVATION



LINEAR TIME-INVARIANT (L)



More generally, for a linear time-variant system:



where $G(t, t_0)$ is the response of the system to an impulse at time t_0 .

$G(t, t_0)$ is called the GREEN'S FUNCTION of the ODE system.

With initial conditions: (when $u(t) = 0$ for $t < 0$)

$$\frac{du}{dt} = \mathcal{L}(u(t)) + f(t) \quad \text{with I.C. } u(0) = u_0$$

Equivalently:

$$\frac{du}{dt} = \mathcal{L}(u(t)) + f(t) + u_0 \delta(t) \quad \text{with zero I.C.}$$

Equivalent because:

$$\left(\begin{array}{l} 0^- = -\epsilon \\ 0^+ = +\epsilon \end{array} \text{ for } \lim_{\epsilon \rightarrow 0} \right)$$

$$\underbrace{\int_{0^-}^{0^+} \frac{du}{dt} dt}_{u(0^+) - u(0^-)} = \underbrace{\int_{0^-}^{0^+} (\mathcal{L}(u(t)) + f(t)) dt}_{=0} + u_0 \underbrace{\int_{0^-}^{0^+} \delta(x) dt}_{1 \text{ by definition}}$$

$$\text{or: } u(0^+) = u(0^-) + u_0 = u_0$$

(i.e., u steps from 0 to u_0 at $t=0$)

$$\Rightarrow u(t) = \int_{-\infty}^t (u_0 \delta(t_0) + f(t_0)) G(t, t_0) dt_0 \quad \text{where } f(t_0) = 0 \text{ for } t_0 < 0$$

$$= u_0 G(t, 0) + \int_0^t f(t_0) G(t, t_0) dt_0$$

↓
I.C.

↓
GREEN'S
FUNCTION
@ $t_0 = 0$

↓
DRIVING
SOURCE

↓
GREEN'S
FUNCTION
@ t_0

GREEN'S FUNCTION FOR PDES ON BOUNDED DOMAINS

1) HOMOGENEOUS PDES (1 dimension)

e.g.: $\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$ with $\begin{cases} u(0, t) = 0 & \text{B.C.} \\ u(L, t) = 0 & \text{B.C.} \\ u(x, 0) = g(x) & \text{I.C.} \end{cases}$

Separation of variables:

$$u(x, t) = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{L} e^{-k \left(\frac{n\pi}{L}\right)^2 t} \quad \text{with} \quad a_n = \frac{2}{L} \int_0^L g(x_0) \sin \frac{n\pi x_0}{L} dx_0$$

Substitution gives:

$$u(x, t) = \int_0^L g(x_0) \left(\underbrace{\sum_{n=1}^{\infty} \frac{2}{L} \sin \frac{n\pi x_0}{L} \sin \frac{n\pi x}{L} e^{-k \left(\frac{n\pi}{L}\right)^2 t}}_{G(x, t; x_0)} \right) dx_0$$

More generally: I.C. $u(x, t_0) = g(x)$ at time t_0

$$\Rightarrow u(x, t) = \int_0^L g(x_0) G(x, t; x_0, t_0) dx_0$$

$$\begin{aligned} \text{with } G(x, t; x_0, t_0) &= \sum_{n=1}^{\infty} \frac{2}{L} \sin \frac{n\pi x_0}{L} \sin \frac{n\pi x}{L} e^{-k \left(\frac{n\pi}{L}\right)^2 (t-t_0)} \\ &= G(x, t-t_0; x_0, 0) \quad (\text{LTI PDE}) \end{aligned}$$

Proof: $\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x,t)$ with $\begin{cases} u(x,0) = g(x) & \text{I.C.} \\ u(0,t) = 0 & \text{B.C. @ 0} \\ u(L,t) = 0 & \text{B.C. @ L} \end{cases}$
 (Sec 9.2 pp 381-383)

a. Homogeneous problem: $Q(x,t) = 0$

Separation of variables: $u(x,t) = \sum_{m=1}^{\infty} \underbrace{a_m}_{\text{COEFFICIENT}} \underbrace{\sin\left(\frac{m\pi x}{L}\right) \cdot e^{-k\left(\frac{m\pi}{L}\right)^2 t}}_{\text{EIGENMODES}}$

b. Inhomogeneous problem: $Q(x,t) \neq 0$

Variation of coefficients: $u(x,t) = \sum_{m=1}^{\infty} a_m(t) \sin\left(\frac{m\pi x}{L}\right) e^{-k\left(\frac{m\pi}{L}\right)^2 t}$

$$\frac{\partial u}{\partial t} = \sum_{m=1}^{\infty} \frac{da_m}{dt} \sin\left(\frac{m\pi x}{L}\right) e^{-k\left(\frac{m\pi}{L}\right)^2 t} + \sum_{m=1}^{\infty} a_m(t) \sin\left(\frac{m\pi x}{L}\right) \left(-k\left(\frac{m\pi}{L}\right)^2\right) e^{-k\left(\frac{m\pi}{L}\right)^2 t}$$

$$= k \frac{\partial^2 u}{\partial x^2} + Q = k \sum_{m=1}^{\infty} a_m(t) \left(-\left(\frac{m\pi}{L}\right)^2\right) \sin\left(\frac{m\pi x}{L}\right) e^{-k\left(\frac{m\pi}{L}\right)^2 t} + Q(x,t)$$

$$\Rightarrow \sum_{m=1}^{\infty} \left(\frac{da_m}{dt} e^{-k\left(\frac{m\pi}{L}\right)^2 t} \right) \sin\left(\frac{m\pi x}{L}\right) = Q(x,t) = \sum_{m=1}^{\infty} q_m(t) \sin\left(\frac{m\pi x}{L}\right)$$

$q_m(t) = \frac{2}{L} \int_0^L Q(x,t) \sin\left(\frac{m\pi x_0}{L}\right) dx_0$
 FOURIER SERIES EXPANSION of $Q(x,t)$

$$\Rightarrow \frac{da_m}{dt} e^{-k\left(\frac{m\pi}{L}\right)^2 t} = q_m(t) \Rightarrow \int_{a_m(0)}^{a_m(t)} da_m = \int_0^t q_m(t_0) e^{+k\left(\frac{m\pi}{L}\right)^2 t_0} dt_0$$

$$\Rightarrow a_m(t) = \underbrace{a_m(0)}_{\text{FOURIER SERIES EXPANSION of } g(x) \text{ I.C.}} + \int_0^t \underbrace{q_m(t_0)}_{\text{FOURIER SERIES EXPANSION of } Q(x,t) \text{ DRIVING SOURCE}} e^{k\left(\frac{m\pi}{L}\right)^2 t_0} dt_0$$

$$\text{I.C. : } g(x) = u(x,0) = \sum_{n=1}^{\infty} d_n(0) \sin\left(\frac{n\pi x}{L}\right)$$

$$\rightarrow d_n(0) = \frac{2}{L} \int_0^L g(x_0) \sin\left(\frac{n\pi x_0}{L}\right) dx_0$$

FOURIER SERIES EXPANSION of $g(x)$ I.C.

$$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} \left(d_n(0) + \int_0^t q_n(t_0) e^{-k\left(\frac{n\pi}{L}\right)^2 t_0} dt_0 \right) \cdot \sin\left(\frac{n\pi x}{L}\right) \cdot e^{-k\left(\frac{n\pi}{L}\right)^2 t}$$

$$\frac{2}{L} \int_0^L g(x_0) \sin\left(\frac{n\pi x_0}{L}\right) dx_0$$

$$\frac{2}{L} \int_0^L Q(x_0, t_0) \sin\left(\frac{n\pi x_0}{L}\right) dx_0$$

Interchanging the order of $\sum_{n=1}^{\infty} \dots$ and $\int_0^L dx_0 \dots$ and $\int_0^t dt_0 \dots$:

$$u(x,t) = \int_0^L g(x_0) \left(\sum_{n=1}^{\infty} \frac{2}{L} \sin\left(\frac{n\pi x_0}{L}\right) \sin\left(\frac{n\pi x}{L}\right) e^{-k\left(\frac{n\pi}{L}\right)^2 t} \right) dx_0$$

I.C. $G(x,t; x_0, 0)$

$$+ \int_0^L \int_0^t Q(x_0, t_0) \left(\sum_{n=1}^{\infty} \frac{2}{L} \sin\left(\frac{n\pi x_0}{L}\right) \sin\left(\frac{n\pi x}{L}\right) e^{-k\left(\frac{n\pi}{L}\right)^2 (t-t_0)} \right) dx_0 dt_0$$

DRIVING SOURCE $G(x,t; x_0, t_0)$

NOTES:

- For LTI (linear time invariant) PDEs:

$$G(x, t; x_0, t_0) = G(x, t - t_0; x_0, 0)$$

- also, for LSI (linear space invariant) PDEs:

$$G(x, t; x_0, t_0) = G(x - x_0, t; 0, t_0)$$

- and thus for LSTI (linear space and time invariant) PDEs:

$$G(x, t; x_0, t_0) = G(x - x_0, t - t_0; 0, 0)$$

- For all linear PDEs:

$$G(x, t; x_0, t) = \delta(x - x_0)$$

because of causality!

- The Green's function for general linear PDEs of the form:

$$\mathcal{L}_{x,t} u(x, t) = f(x, t) \quad \text{with} \quad \begin{cases} \text{B.C.} = h(t) \\ \text{I.C.} = g(x) \end{cases}$$

can be found as the solution to:

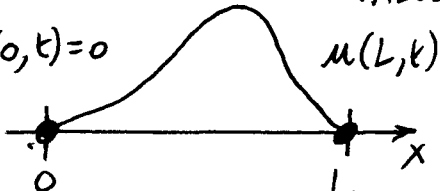
$$\mathcal{L}_{x,t} G(x, t; x_0, t_0) = \delta(x - x_0) \delta(t - t_0) \quad \text{with} \quad \begin{cases} \text{B.C.} = 0 \\ \text{I.C.} = 0 \end{cases}$$

or, equivalently:

$$\mathcal{L}_{x,t} G(x, t; x_0, t_0) = 0 \quad \text{with} \quad \begin{cases} \text{B.C.} = 0 \\ \text{I.C.} = \delta(x - x_0) \end{cases}$$

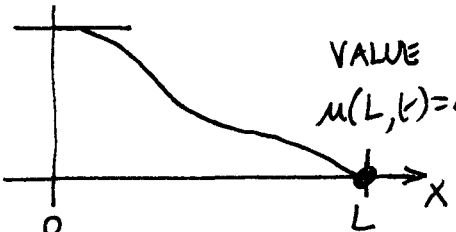
• The Green's function depends on the type of homogeneous B.C., e.g.:

VALUE $u(0,t)=0$ VALUE $u(L,t)=0$



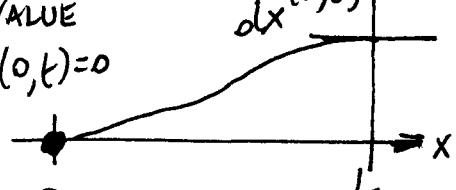
$G(x,t;x_0,t_0) = \sum_{n=1}^{\infty} \frac{2}{L} \sin \frac{n\pi x_0}{L} \sin \frac{n\pi x}{L} e^{-k \left(\frac{n\pi}{L}\right)^2 (t-t_0)}$

FLUX $\frac{du}{dx}(0,t)=0$ VALUE $u(L,t)=0$



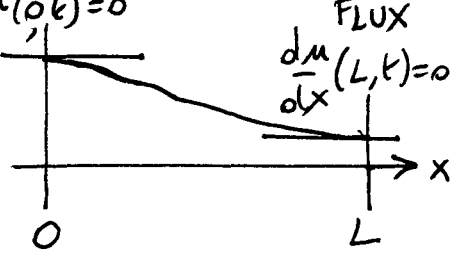
$G(x,t;x_0,t_0) = \sum_{n=\text{odd}}^{\infty} \frac{2}{L} \cos \frac{n\pi x_0}{2L} \cos \frac{n\pi x}{2L} e^{-k \left(\frac{n\pi}{2L}\right)^2 (t-t_0)}$
($n=1,3,5,\dots$)

VALUE $u(0,t)=0$ FLUX $\frac{du}{dx}(L,t)=0$



$G(x,t;x_0,t_0) = \sum_{n=\text{odd}}^{\infty} \frac{2}{L} \sin \frac{n\pi x_0}{2L} \sin \frac{n\pi x}{2L} e^{-k \left(\frac{n\pi}{2L}\right)^2 (t-t_0)}$
($n=1,3,5,\dots$)

FLUX $\frac{du}{dx}(0,t)=0$ FLUX $\frac{du}{dx}(L,t)=0$



$G(x,t;x_0,t_0) = \sum_{n=1}^{\infty} \frac{2}{L} \cos \frac{n\pi x_0}{L} \cos \frac{n\pi x}{L} e^{-k \left(\frac{n\pi}{L}\right)^2 (t-t_0)}$
 $+ \frac{1}{L}$

Physical interpretation:

"VALUE" $u(x,t)$: TEMPERATURE ; VOLTAGE \rightarrow "SHORT"

"FLUX" $k \frac{du}{dx}(x,t)$: HEAT TRANSFER ; CURRENT \rightarrow "OPEN CIRCUIT"