

Lecture 16

Finite-Element Methods

References

Haberman APDE, Sec. 6.7.

http://en.wikipedia.org/wiki/Finite_element_method

FINITE ELEMENT METHODS

(Haberman Sec. 6.7)

- Separation of variables ^(in space) works for simple problems, with geometries that allow for separation of variables in the boundary conditions.
- Finite element methods offer numerical solutions to diffusion, wave, and other linear PDE problems in bioengineering with complex, arbitrary boundary conditions on real-life geometries (not cubes, cylinders, or spheres).

General linear PDEs (in homogeneous media):

$$L_t(u(\vec{r}, t)) = \Delta u(\vec{r}, t) + f(\vec{r}, t)$$

e.g. {

DIFFUSION EQ.: $L_t(u) = \frac{1}{D} \frac{\partial u}{\partial t}$ $[D: \frac{m^2}{s}]$

WAVE EQ.: $L_t(u) = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$ $[c: \frac{m}{s}]$

CABLE EQ.: $L_t(u) = lC \frac{\partial^2 u}{\partial t^2} + rC \frac{\partial u}{\partial t}$

$[l: \frac{H}{m}]$ $[C: \frac{F}{m}]$ $[r: \frac{\Omega}{m}]$

General (inhomogeneous) problem:

$$\mathcal{L}_t(\mu) = \Delta \mu + \underbrace{f(\vec{r}, t)}_{\text{SOURCE}}$$

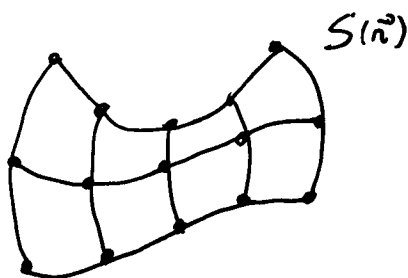
with: I.C. : $\mu(\vec{r}, 0) = g(\vec{r})$, and $\frac{\partial \mu}{\partial t}(\vec{r}, 0) = h(\vec{r})$
DIFF/WAVE/CABLE WAVE/CABLE

B.C. : • VALUE B.C. : $\mu(S(\vec{r}), t) = \mu_V(S(\vec{r}))$

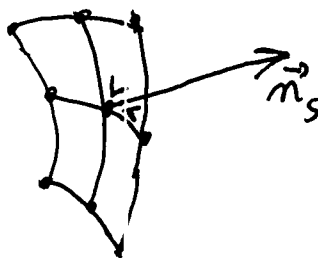
AND/OR

• FLUX B.C. : $\vec{\nabla} \mu(s, t) \cdot \vec{n}_S = \mu_F(S(\vec{r}))$

where $S(\vec{r})$ defines a SURFACE in space



and \vec{n}_S is locally perpendicular to the surface S .



Homogeneous problem:

$$\mathcal{L}_t(\mu_H) = \Delta \mu_H$$

$$\text{with } \begin{cases} \text{SOURCE} = 0 & f(\vec{r}, t) = 0 \\ \text{VALUE B.C.} = 0 & \mu_H(S(\vec{r}), t) = 0 \\ \text{AND/OR} & \\ \text{FLUX B.C.} = 0 & \vec{\nabla} \mu_H(S, t) \cdot \vec{n}_S = 0 \end{cases}$$

Separation of space and time (but NOT separation in space):

$$\mu_H(\vec{r}, t) = \phi(\vec{r}) \cdot G(t)$$

TIME: $\mathcal{L}_t(G(t)) = -\lambda G(t)$

• DIFF.: $\mathcal{L}_t(G) = \frac{1}{D} \frac{d}{dt} G \Rightarrow G(t) = G(0) \cdot e^{-D, \lambda t}$

• WAVE: $\mathcal{L}_t(G) = \frac{1}{c^2} \frac{d^2}{dt^2} G \Rightarrow$

$$G(t) = G(0) \cdot \cos(\sqrt{\lambda} ct)$$

$$+ \frac{\dot{G}(0)}{\sqrt{\lambda} c} \cdot \sin(\sqrt{\lambda} ct)$$

• CABLE: $\mathcal{L}_t(G) = lC \frac{d^2 G}{dt^2} + rC \frac{dG}{dt}$

(left as an exercise)

SPACE : $\Delta \phi(\vec{r}) + \lambda \phi(\vec{r}) = 0$ (Haberman p.270)

with complex B.C. (value/flux)

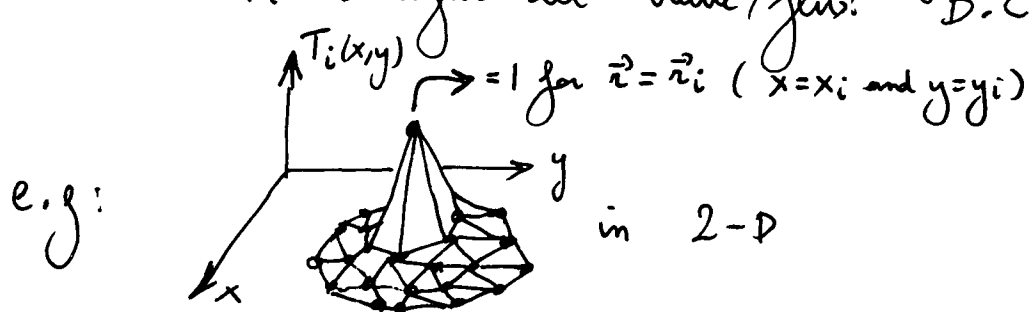
→ Solve using FINITE ELEMENTS

- specific to the B.C. geometry and type (value/flux)

- independent of PDE type (diffusion; wave; cable; ...)

Let $\phi(\vec{r}) = \sum_{i=1}^n U_i \cdot T_i(\vec{r})$ on a grid with discrete points \vec{r}_i such that:

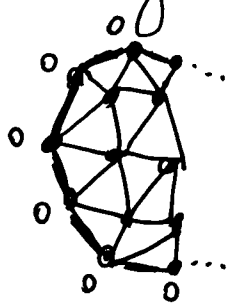
- $T_i(\vec{r}_i) = 1$
- $T_i(\vec{r}_j) = 0$ for $j \neq i$
- T_i is PIECEWISE LINEAR on segments connecting vertices on the grid
- T_i satisfies all value/flux B.C.



How does T_i satisfy the boundary conditions?

• VALUE B.C. : $T_i(S(i)) = 0$

\Rightarrow T_i is zero, for each of the i , on each of the vertex points that are on the boundary (as well as the segments that connect them).



• FLUX B.C. : $\vec{\nabla} T_i(s) \cdot \vec{n}_s = 0$

\Rightarrow The gradient (slope) of T_i , for each of the i , is PARALLEL to the surface segments suspended by the vertex points on the boundary.

\Rightarrow The average of T_i values on the vertex points of any polyhedron segment on the boundary should equal the T_i value of the vertex interior point, part of the same polyhedron.

We seek a solution $\phi(\vec{r}) = \sum_{j=1}^n u_j T_j(\vec{r})$ that approximates

$$\Delta \phi(\vec{r}) + \lambda \phi(\vec{r}) = 0$$

in integral form:

$$\iiint_V T_i(\vec{r}) \cdot (\Delta \phi(\vec{r}) + \lambda \phi(\vec{r})) d^3r = 0, \quad i=1,2,\dots,n$$

over the domain (volume) V enclosed by the B.C. surface S .

Note: $\vec{\nabla} \cdot (T_i \vec{\nabla} \phi) = T_i \cdot \underbrace{\vec{\nabla} \cdot \vec{\nabla} \phi}_{\Delta \phi} + \vec{\nabla} T_i \cdot \vec{\nabla} \phi$

$$\Rightarrow \iiint_V T_i \Delta \phi d^3r = \underbrace{\iiint_V \vec{\nabla} \cdot (T_i \vec{\nabla} \phi) d^3r}_{\text{GAUSS}} - \iiint_V \vec{\nabla} T_i \cdot \vec{\nabla} \phi d^3r$$

$$= \oint_{S(V)} T_i \vec{\nabla} \phi \cdot \vec{n} dS = 0 \text{ because:}$$

- $T_i = 0$ ON VALUE BOUNDARY
- $\vec{\nabla} \phi \cdot \vec{n} = 0$ ON FLUX BOUNDARY
($\vec{\nabla} T_j \cdot \vec{n} = 0$)

$$\Rightarrow \iiint_V \vec{\nabla} T_i \cdot \vec{\nabla} \phi d^3r = \lambda \iiint_V T_i \cdot \phi d^3r$$

GALERKIN APPROXIMATION

(WEAK FORM)

(Haberman pp. 267-270)

$$\Rightarrow \sum_{j=1}^n \underbrace{\iiint_V \vec{\nabla} T_i \cdot \vec{\nabla} T_j d^3\vec{r}}_{K_{ij}} \cdot u_j = \lambda \sum_{j=1}^n \underbrace{\iiint_V T_i T_j d^3\vec{r}}_{M_{ij}} \cdot u_j$$

STIFFNESS MATRIX \bar{K} MASS MATRIX \bar{M}

$$\Rightarrow \bar{K} \vec{u} = \lambda \bar{M} \vec{u}$$

\Rightarrow generalized eigenvector-eigenvalue problem for non-zero solutions \vec{u} :

$$\det(\bar{K} - \lambda \bar{M}) = 0 \quad \text{with } \vec{u} \text{ in the null-space of } \bar{K} - \lambda \bar{M}$$

Solutions to λ and \vec{u} are eigenvalues and eigenvectors of $\bar{M}^{-1} \cdot \bar{K}$

For any eigenvalue λ , the solution $\phi(\vec{r}) = \sum_{j=1}^n u_j T_j(\vec{r})$ is the corresponding EIGENMODE.

Since $\bar{M}^{-1} \cdot \bar{K}$ is $(n \times n)$, there are n eigenvalues λ_k and n corresponding eigenmodes $\phi_k(\vec{r}) = \sum_{j=1}^n (u_k)_j T_j(\vec{r})$

$$\Rightarrow u_H(\vec{r}, t) = \sum_k A_k \phi_k(\vec{r}) \cdot G_k(t)$$

HOMOGENEOUS SOLUTION

with $G_k(t) = e^{-D\lambda_k t}$ for DIFFUSION.

How to find the coefficients A_k ?

As with all diffusion/wave problems (satisfying Sturm-Liouville),
the eigenmodes $\phi_k(\vec{r})$ are ORTHOGONAL:

Proof: \bar{M} and \bar{K} are symmetric (and positive definite)

\Rightarrow eigenvalues λ_k are real (and positive),

and eigenvectors \vec{u}_k are orthogonal w.r.t. \bar{M} :

$$\vec{u}_k^T \cdot \bar{M} \cdot \vec{u}_l = c_k \delta_{kl}$$

$$\begin{aligned} \Rightarrow \iiint_V \phi_k(\vec{r}) \phi_l(\vec{r}) d^3\vec{r} &= \sum_{i,j} (u_k)_i (u_l)_j \underbrace{\iiint_V T_i(\vec{r}) T_j(\vec{r}) d^3\vec{r}}_{M_{ij}} \\ &= \vec{u}_k^T \cdot \bar{M} \cdot \vec{u}_l = c_k \delta_{kl} \quad \text{Q.E.D.} \end{aligned}$$

$$\begin{aligned} \text{I.C. } u_H(\vec{r}, 0) &= \sum_k A_k \phi_k(\vec{r}) \cdot \underset{\substack{\text{"} \\ \text{1 for DIFFUSION}}}{G_k(0)} \\ &= g(\vec{r}) \end{aligned}$$

$$\Rightarrow A_k = \frac{\iiint_V g(\vec{r}) \phi_k(\vec{r}) d^3\vec{r}}{\iiint_V [\phi_k(\vec{r})]^2 d^3\vec{r}} \quad (\text{as before})$$

General solution to INHOMOGENEOUS PDE with INHOMOGENEOUS B.C.:

use Green's functions as before.

The Green's function is the homogeneous solution $u_H(\vec{r}, t)$
with I.C. $u_H(\vec{r}, t') = \delta(\vec{r} - \vec{r}')$

↓
INITIAL
ACTIVATION
AT TIME t'

↓
IMPULSE
AT LOCATION \vec{r}'

for DIFFUSION

$$\Rightarrow \sum_k A_k \phi_k(\vec{r}) \cdot e^{-D\lambda_k t'} = \delta(\vec{r} - \vec{r}')$$

$$\Rightarrow \sum_k A_k \underbrace{\iiint_V \phi_k(\vec{r}) \phi_l(\vec{r}) d^3\vec{r}}_{c_k \delta_{kl}} e^{-D\lambda_k t'} = \underbrace{\iiint_V \delta(\vec{r} - \vec{r}') \phi_l(\vec{r}) d^3\vec{r}}_{\phi_l(\vec{r}')}$$

$$\Rightarrow A_l = \frac{1}{c_l} \phi_l(\vec{r}') \cdot e^{+D\lambda_l t'}$$

$$G(\vec{r}, t; \vec{r}', t') = u_H(\vec{r}, t) = \sum_k A_k \phi_k(\vec{r}) \cdot e^{-D\lambda_k t}$$

$$= \sum_k \frac{1}{c_k} \cdot \phi_k(\vec{r}') \phi_k(\vec{r}) e^{-D\lambda_k(t-t')}$$

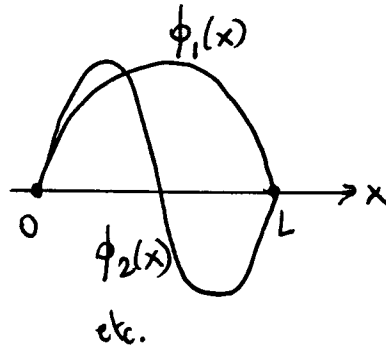
(as before!)

$$\text{with } c_k = \iiint_V (\phi_k(\vec{r}'))^2 d^3\vec{r}$$

EXAMPLE : Diffusion in 1-D

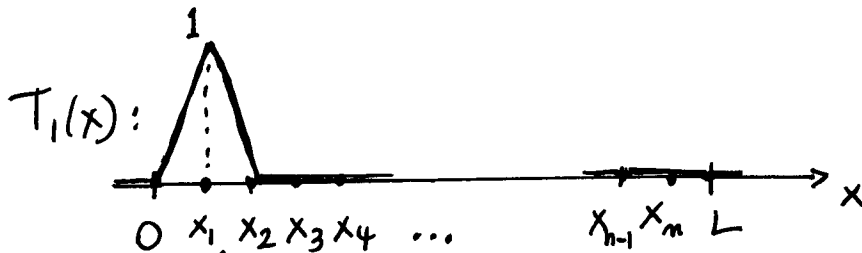
on $[0, L]$ with VALUE-VALUE B.C.

$$\Rightarrow u_H(x, t) = \sum_k A_k \underbrace{\sin\left(\frac{k\pi x}{L}\right)}_{\phi_k(x)} \cdot \underbrace{e^{-D\left(\frac{k\pi}{L}\right)^2 t}}_{G_k(t)}$$

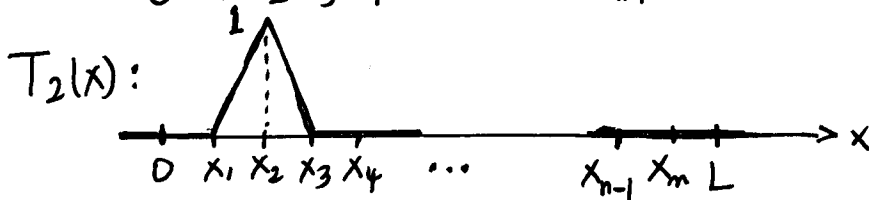


$$\lambda_k = \left(\frac{k\pi}{L}\right)^2$$

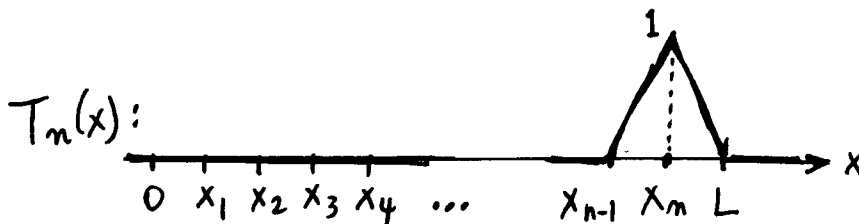
ANALYTICAL SOLUTION



FINITE ELEMENTS



⋮
↔
 Δx



$$\Rightarrow K_{ij} = \int_0^L \frac{dT_i}{dx}(x) \frac{dT_j}{dx}(x) dx = \begin{cases} \frac{2}{\Delta x} & \text{for } i=j \\ -\frac{1}{\Delta x} & \text{for } i=j \pm 1 \end{cases}$$

$$M_{ij} = \int_0^L T_i(x) T_j(x) dx = \begin{cases} \frac{2}{3} \Delta x & \text{for } i=j \\ \frac{1}{6} \Delta x & \text{for } i=j \pm 1 \end{cases}$$

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%% Finite Elements: Linear (1-D) Diffusion with value-value B.C.
%% BENG 221 example, 11/17/2011

L = 1; % length
n = 100; % number of finite elements
ne = 10; % number of displayed eigenmodes
don = diag(ones(n,1),0); % on-diagonal, all ones
doff = diag(ones(n-1,1),1); % off-diagonal, all ones

dx = 1 / (n+1); % grid spacing, uniform
x = (0:n+1)' * L / (n+1); % vertex points, including 0 and L boundary points
K = 1/dx * (2*don - doff - doff'); % stiffness matrix
M = dx/6 * (4*don + doff + doff'); % mass matrix

% finite differences: eigenvalue-eigenmode problem K*U = Lambda*dx*U with Lambda diagonal
help eig
% [U, Lambda] = eig(K/dx);
[U, Lambda] = eig(K, dx*don);
lambda_fd = diag(Lambda); % eigenvalues
phi_fd = [zeros(1,n); U; zeros(1,n)]; % eigenmodes; including 0 and L boundary points
phi_fd = phi_fd * diag(sign(phi_fd(2,:))); % flip eigenmodes with negative slope at zero

% finite elements: generalized eigenvalue-eigenmode problem K*U = Lambda*M*U with Lambda diagonal
[U, Lambda] = eig(K, M);
lambda_fe = diag(Lambda); % eigenvalues
phi_fe = [zeros(1,n); U; zeros(1,n)]; % eigenmodes; including 0 and L boundary points
phi_fe = phi_fe * diag(sign(phi_fe(2,:))); % flip eigenmodes with negative slope at zero

% analytical solution; ground truth
lambda_anal = (pi * (1:n) / L).^2; % analytical eigenvalues (separation of variables)
phi_anal = sqrt(2) * sin(pi * x * (1:n) / L); % analytical eigenmodes (separation of variables)

% plot eigenvalues
figure(1)
semilogy(lambda_fd, 'k:')
hold on
semilogy(lambda_fe, 'k-')
semilogy(lambda_anal, 'k--')
hold off
xlabel('Mode{\it k}')
ylabel('Eigenvalue {\it \lambda_k}')
legend('Finite Differences', 'Finite Elements', 'Analytical', 'Location', 'NorthWest')
print -dpdf 'linfinelem_eigenvalues.pdf'

% plot eigenmodes
figure(2)
plot(x, phi_fd(:,1:ne), ':')
hold on
plot(x, phi_fe(:,1:ne), '-')
plot(x, phi_anal(:,1:ne), '--')
hold off
xlabel('{\it x/L}')
ylabel('Eigenmodes {\it \phi_k}')
print -dpdf 'linfinelem_eigenmodes.pdf'

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