

## Lecture 3

# Canonical LTI ODEs, Eigenmode Analysis, and Principal Component Analysis

### References

[http://en.wikipedia.org/wiki/LTI\\_system\\_theory](http://en.wikipedia.org/wiki/LTI_system_theory)

[https://en.wikipedia.org/wiki/Eigenvalues\\_and\\_eigenvectors](https://en.wikipedia.org/wiki/Eigenvalues_and_eigenvectors)

[https://en.wikipedia.org/wiki/Principal\\_component\\_analysis](https://en.wikipedia.org/wiki/Principal_component_analysis)

[https://en.wikipedia.org/wiki/Singular-value\\_decomposition](https://en.wikipedia.org/wiki/Singular-value_decomposition)

# EIGEN ANALYSIS OF CANONICAL LTI ODES:

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Canonical form:

$$\frac{d\vec{u}}{dt} = \underset{\text{MATRIX}}{\bar{A}} \cdot \underset{\text{VECTOR}}{\vec{u}} + \vec{f} \quad \text{with} \quad \vec{u}(0) = \vec{u}_0$$

I.C.

Homogeneous case:  $\vec{f} = 0$

$$\frac{d\vec{u}}{dt} = \bar{A} \cdot \vec{u}$$

Trivial solution:  $\vec{u}(t) \equiv 0$

Non-trivial solutions are the EIGENMODES  $\vec{u}_i(t) \neq 0$

where in general  $\vec{u}(t) = \sum_i c_i \vec{u}_i(t)$

These eigenmodes can be found by  
SEPARATION OF VARIABLES in terms of  
EIGENVECTORS and EIGENVALUES of  $\bar{A}$ .

The general solution can also be found by Laplace  
in terms of the same eigen decomposition.

## SEPARATION OF VARIABLES :

Try a non-trivial solution to the homogeneous ODE :

$$\frac{d\vec{u}}{dt} = \bar{A} \cdot \vec{u}$$

of the form, separated in time and (vector) space :

$$\vec{u}(t) = \vec{u} \cdot b(t)$$

↓  
VECTOR  
(SPACE)  
independent of  
TIME

↓  
SCALAR  
function of  
TIME

$$\Rightarrow \frac{d}{dt} (\vec{u} b(t)) = \vec{u} \frac{db}{dt} = \bar{A} \cdot \vec{u} b(t)$$

This identity can only non-trivially hold when  $\bar{A} \cdot \vec{u}$   
and  $\vec{u}$  are co-linear :

$$\bar{A} \cdot \vec{u} = \lambda \vec{u} \quad \lambda \neq 0 \text{ (NON-TRIVIAL)}$$

i.e.,  $\vec{u}$  is an EIGENVECTOR of  $\bar{A}$  with  
EIGENVALUE  $\lambda$ .

$$\Rightarrow \vec{u} \frac{d\psi}{dt} = \lambda \vec{u} \psi(t) \quad \text{where } \|\vec{u}\| \neq 0 \text{ (NON-TRIVIAL)}$$

$$\Rightarrow \frac{d\psi}{dt} = \lambda \psi \quad \text{or} \quad d(\ln \psi) = \lambda dt$$

$$\ln \psi = \lambda t + c t$$

$$\psi(t) = e^{\lambda t} \cdot \underbrace{e^{c t}}_{\text{another constant, } c}$$

$$\psi(t) = c \cdot e^{\lambda t}$$

$$\Rightarrow \vec{u}(t) = \vec{u} e^{\lambda t} \text{ is ONE solution to the homogeneous ODE}$$

where  $\vec{A} \cdot \vec{u} = \lambda \vec{u}$ .

In general there are  $N$  such eigenvectors/eigenvalues:

$$\vec{A} \cdot \vec{u}_i = \lambda_i \vec{u}_i \quad i=1, \dots, N$$

each contributing such solution (EIGENMODE):

$$\vec{u}_i(t) = \vec{u}_i \cdot e^{\lambda_i t}$$

where  $N$  is the dimension of  $\vec{u}$  (ORDER of the ODE).

The general solution is thus a superposition of EIGENMODES:

$$\vec{u}(t) = \sum_{i=1}^N c_i \vec{u}_i(t) = \sum_{i=1}^N c_i \vec{u}_i e^{\lambda_i t}$$

The particular solution (values of  $c_i$ ) is determined by the INITIAL CONDITIONS (I.C.s):

$$\vec{u}_0 = \vec{u}(0) = \sum_{i=1}^N c_i \vec{u}_i(0) = \sum_{i=1}^N c_i \vec{u}_i$$

→  $N$  linear equations in the  $N$  unknown  $c_i$ 's

SPECIAL CASE: for  $\bar{A}$  symmetric:

→ the eigenvalues  $\lambda_i$  are REAL

→ the eigenvectors  $\vec{u}_i$  are ORTHOGONAL:  
(ORTHONORMAL upon rescaling)

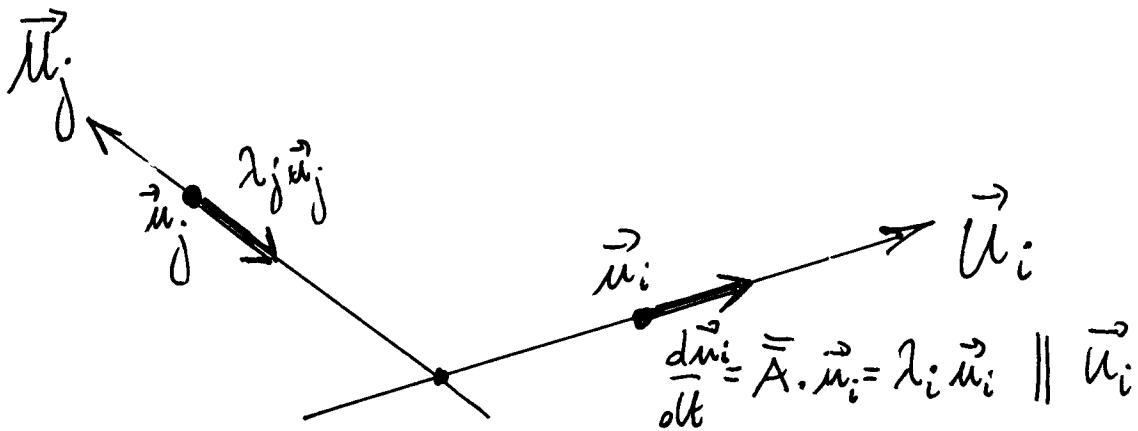
$$\vec{u}_i \cdot \vec{u}_j = \delta_{ij} = \begin{cases} 1 & \text{for } i=j \\ 0 & \text{otherwise} \end{cases}$$

$$\Rightarrow c_i = \vec{u}_0 \cdot \vec{u}_i$$

$$\Rightarrow \vec{u}(t) = \sum_{i=1}^N (\vec{u}_0 \cdot \vec{u}_i) \vec{u}_i e^{\lambda_i t}$$

# EIGENVECTORS AS PRINCIPAL AXES:

The eigenvectors of  $\bar{A}$  are those directions in state space  $\vec{u}$  for which the homogeneous dynamics is decoupled and self-contained:



Eigen decomposition and transformation to principal axes:

$$\bar{A} = \bar{U} \cdot \bar{\Lambda} \cdot \bar{U}^{-1}$$

with  $\bar{U} = (\vec{u}_1 \vec{u}_2 \dots \vec{u}_m)$  eigenvector matrix

$$\bar{\Lambda} = \begin{pmatrix} \lambda_1 & & & 0 \\ & \lambda_2 & & \\ & & \ddots & \\ 0 & & & \lambda_m \end{pmatrix} \text{ diagonal eigenvalue matrix}$$

Equivalently:  $\bar{A} \cdot \bar{U} = \bar{U} \cdot \bar{\Lambda}$  or  $\bar{A} \cdot \vec{u}_i = \lambda_i \vec{u}_i$

$$\begin{aligned} \frac{d\vec{u}}{dt} &= \vec{A} \cdot \vec{u} + \vec{f} \\ &= \vec{U} \cdot \vec{\Lambda} \cdot \vec{U}^{-1} \cdot \vec{u} + \vec{f} \end{aligned}$$

$$\Downarrow$$

$$\vec{U}^{-1} \frac{d\vec{u}}{dt} = \cancel{\vec{U}^{-1} \vec{U}} \cdot \vec{\Lambda} \cdot \vec{U}^{-1} \vec{u} + \vec{U}^{-1} \vec{f}$$

$$\Downarrow$$

$$\frac{d\vec{u}'}{dt} = \vec{\Lambda} \cdot \vec{u}' + \vec{f}'$$

where

$$\begin{cases} \vec{u}' = \vec{U}^{-1} \cdot \vec{u} \\ \vec{f}' = \vec{U}^{-1} \cdot \vec{f} \end{cases} \quad \begin{array}{l} \text{projections onto} \\ \text{principal axes } \vec{U} \end{array}$$

$$\frac{du_i'}{dt} = \lambda_i u_i' + f_i'$$

$$u_i'(t) = u_i'(0) e^{\lambda_i t} + \int_0^t f_i'(\theta) e^{\lambda_i(t-\theta)} d\theta$$

I.c. ;

DRIVING FORCE :

projected on principal axis  $U_i$

$$\text{and } \vec{u}(t) = \sum_i u_i'(t) \cdot \vec{U}_i$$

$$\Rightarrow \vec{u}(t) = \sum_i u_i^?(0) e^{\lambda_i t} \vec{u}_i + \sum_i \int_0^t f_i^?(\theta) e^{\lambda_i(t-\theta)} d\theta \cdot \vec{u}_i$$

where :

$$\begin{cases} u_i^?(0) = (\bar{U}^{-1} \vec{u}(0))_i : & \text{PROJECTION OF I.C. } \vec{u}(0) \\ & \text{ONTO PRINCIPAL AXIS } i \\ f_i^?(0) = (\bar{U}^{-1} \vec{f}(0))_i : & \text{PROJECTION OF STIMULUS } \vec{f}(0) \\ & \text{ONTO PRINCIPAL AXIS } i \end{cases}$$

Special case:  $\bar{A} = \text{symmetric}$

$$\Rightarrow \bar{U}^{-1} = \bar{U}^T \quad (\text{and } \lambda_i = \text{real})$$

Then:

$$\begin{cases} u_i^?(0) = \vec{u}_i^T \vec{u}(0) = \vec{u}_i \cdot \vec{u}(0) \\ f_i^?(0) = \vec{u}_i^T \vec{f}(0) = \vec{u}_i \cdot \vec{f}(0) \end{cases} \quad \text{DOT PRODUCTS}$$

And thus:

$$\vec{u}(t) = \sum_i (\vec{u}_i \cdot \vec{u}(0)) e^{\lambda_i t} \vec{u}_i + \sum_i \int_0^t (\vec{u}_i \cdot \vec{f}(\theta)) e^{\lambda_i(t-\theta)} d\theta \cdot \vec{u}_i$$



## NOTES:

- 1) The characteristic equation  $\det(\bar{A} - \lambda \bar{I}) = 0$  to find the eigenvalues  $\lambda_1, \lambda_2, \dots, \lambda_m$  is the SAME polynomial zero equality (denominator of Laplace transfer function) to find the poles  $p_1, p_2, \dots, p_m$ , where  $s = \lambda$ :

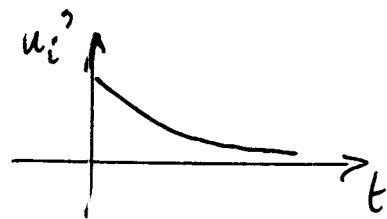
$$s \cdot \vec{U}(s) - \vec{u}(0) = \bar{A} \cdot \vec{U}(s) + \vec{F}(s)$$

$$(s \bar{I} - \bar{A}) \cdot \vec{U}(s) = \vec{u}(0) + \vec{F}(s)$$

- 2) The dynamics is STABLE when all eigenvalues have negative real components:

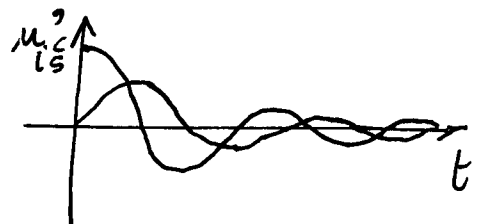
$$\operatorname{Re}(\lambda_i) \leq 0 \quad \forall i$$

— REAL:  $\lambda_i \leq 0$



— COMPLEX CONJUGATE PAIR:

$$\lambda_{i\pm} = -\sigma \pm j\omega$$



3) The eigenvalues with the LARGEST  $\text{Re}(\lambda_i)$  dominate the dynamics  $\vec{u}(t)$  in the long run ( $t \rightarrow \infty$ ):

$$\vec{u} = \sum_i u_i'(t) \cdot \vec{u}_i$$

$$= \sum_i \left( \underbrace{u_i'(0)}_{\text{I.c.}} e^{\lambda_i t} + \int_0^t \underbrace{f_i'(\theta)}_{\substack{\text{DRIVING} \\ \text{FORCE} \\ \text{(finite)}}} e^{\lambda_i(t-\theta)} d\theta \right) \cdot \vec{u}_i$$

$$\propto e^{\text{Re}(\lambda_i) \cdot t} \cdot \vec{u}_i \quad \text{as } t \rightarrow \infty$$

for the largest  $\text{Re}(\lambda_i)$   
 ( $< 0$  if  $\vec{u}$  is stable)

# SINGULAR VALUE DECOMPOSITION (SVD) & PRINCIPAL COMPONENT ANALYSIS (PCA)

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— SVD is an extension to eigen decomposition:

$$\begin{array}{cccc} \overline{X} & = & \overline{U} \cdot \overline{\Sigma} & \cdot \overline{W}^* \\ (n \times p) & & (n \times n) & (n \times p) & (p \times p) \end{array}$$

•  $\overline{X} = (\vec{x}[1], \vec{x}[2], \dots, \vec{x}[p])$   
p points of process  $\vec{x}$  in  
n-dimensional space

•  $\overline{U}$  and  $\overline{W}$  are orthogonal matrices

$$\begin{array}{l} \overline{U}^* \cdot \overline{U} = \overline{I} \quad (n \times n) \\ \overline{W}^* \cdot \overline{W} = \overline{I} \quad (p \times p) \end{array}$$

with \* the complex conjugate transpose

$$(\overline{W}^* = \overline{W}^T \text{ when } \overline{X} \text{ is real})$$

$$\bullet \quad \underline{\underline{\Sigma}} = \begin{pmatrix} \Sigma_1 & 0 & \dots & 0 & 0 & \dots & 0 \\ 0 & \Sigma_2 & & & & & \\ & & \ddots & & & & \\ 0 & 0 & \dots & \Sigma_m & 0 & \dots & 0 \end{pmatrix} \quad (n \times p) \text{ diagonal matrix}$$

with singular values  $\Sigma_1, \Sigma_2, \dots, \Sigma_m$   
(real & positive)

— SVD is an efficient method for PCA :

$\underline{\underline{U}} = (\vec{u}_1, \vec{u}_2, \dots, \vec{u}_m)$  are the principal axes of the data  $\vec{x}[j]$  along which the covariance is diagonal :

$$\begin{aligned} \text{cov}(\vec{x}) &= \frac{1}{p} \sum_{j=1}^p \vec{x}[j] \cdot \vec{x}[j]^* \\ &= \frac{1}{p} \underline{\underline{X}} \cdot \underline{\underline{X}}^* \\ &= \frac{1}{p} (\underline{\underline{U}} \cdot \underline{\underline{\Sigma}} \cdot \underline{\underline{W}}^*) \cdot (\underline{\underline{W}} \cdot \underline{\underline{\Sigma}}^* \cdot \underline{\underline{U}}^*) \\ &= \frac{1}{p} \underline{\underline{U}} \cdot (\underline{\underline{\Sigma}} \cdot \underline{\underline{\Sigma}}^*) \cdot \underline{\underline{U}}^* = \underline{\underline{U}} \cdot \underline{\underline{C}} \cdot \underline{\underline{U}}^* \end{aligned}$$

with diagonal covariance  $\underline{\underline{C}} = \frac{1}{p} \underline{\underline{\Sigma}} \cdot \underline{\underline{\Sigma}}^*$

$$C_i = \frac{1}{p} (\Sigma_i)^2$$

The projections of  $\vec{x} [j]$  onto the principal axes  $\vec{u}_i$  are the PRINCIPAL COMPONENTS  $T_i [j]$ :

$$T_i [j] = \vec{u}_i \cdot \vec{x} [j] \quad \begin{array}{l} i = 1, \dots, n \\ j = 1, \dots, p \end{array}$$

$$\text{or } \bar{\bar{T}} = \bar{U}^* \cdot \bar{X} = \bar{\Sigma} \cdot \bar{W}^*$$

$$\text{with } \bar{\bar{T}} = ( \vec{T} [1], \vec{T} [2], \dots, \vec{T} [p] )$$

Each principal component  $T_i [j]$  explains/captures a portion  $C_i$  of the total variance<sup>(1)</sup>  $\sum_{k=1}^p C_k$  in the data.

Typically, retaining only the PCs with largest  $C_i$  captures most of the variance.

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$$(1) \text{ var}(\vec{x}) = \text{Trace}(\text{cov}(\vec{x})) = \text{Trace}(\bar{C}) = \sum_{k=1}^p C_k$$

## APPENDIX :

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Solution for inhomogeneous LTI ODEs :

$$\frac{d\vec{u}}{dt} = \bar{A} \cdot \vec{u} + \vec{f} \quad \text{with} \quad \vec{u}(0) = \vec{u}_0$$

DRIVING FORCE

I.C.

by :

- 1) Variation of coefficients in the eigenmode expansion
- 2) Laplace and the eigen decomposition of  $\bar{A}$

# INHOMOGENEOUS LTI ODES :

$$\frac{d\vec{u}}{dt} = \vec{A} \cdot \vec{u} + \vec{f} \quad \text{with } \vec{u}(0) = \vec{u}_0$$

$\vec{f}$   
 DRIVING  
 FORCE  
 (SOURCE)

I.C.

- Solution by variation of coefficients in the eigenmode expansion:

$$\text{Let } \vec{u}(t) = \sum_{i=1}^N c_i(t) \vec{u}_i(t) \quad \text{with } \frac{d\vec{u}_i}{dt} = \vec{A} \cdot \vec{u}_i = \lambda_i \cdot \vec{u}_i$$

$c_i(t)$   
 ↓  
 COEFFICIENTS ARE  
 FUNCTION OF TIME

$$\Rightarrow \frac{d\vec{u}}{dt} = \sum_{i=1}^N \frac{dc_i}{dt} \cdot \vec{u}_i + \sum_{i=1}^N c_i \frac{d\vec{u}_i}{dt}$$

$$= \sum_{i=1}^N c_i(t) \cdot \lambda_i \vec{u}_i + \vec{f}$$

$$\text{Let } \vec{f}(t) = \sum_{i=1}^N f_i(t) \vec{u}_i \quad \text{decomposition of } \vec{f} \text{ in basis } \{\vec{u}_i\}$$

$$\text{Also } \vec{u}_i = \vec{u}_i e^{\lambda_i t}$$

$$\Rightarrow \frac{dc_i}{dt} \cdot e^{\lambda_i t} = f_i(t) \quad i = 1, \dots, N$$

$$dc_i = f_i(t) e^{-\lambda_i t} dt$$

$$c_i(t) = c_i(0) + \int_0^t f_i(\theta) e^{-\lambda_i \theta} d\theta$$

$$\Rightarrow \vec{m}(t) = \sum_{i=1}^N c_i(t) \vec{m}_i(t) = \sum_{i=1}^N c_i(t) \vec{u}_i e^{\lambda_i t}$$

$$= \sum_{i=1}^N c_i(0) \vec{u}_i e^{\lambda_i t} + \sum_{i=1}^N \int_0^t f_i(\theta) \vec{u}_i e^{\lambda_i(t-\theta)} d\theta$$

I.C.
ALONG EIGENMODES
DRIVING FORCE
CONVOLVED ALONG EIGENMODES

where  $\vec{m}_0 = \sum_{i=1}^N c_i(0) \vec{u}_i$  I.C. DECOMPOSITION

$\vec{f}(t) = \sum_{i=1}^N f_i(t) \vec{u}_i$  DRIVING FORCE DECOMPOSITION

SPECIAL CASE

$\vec{A}$  = SYMMETRIC:  
 ( $\vec{u}_i$  ORTHONORMAL)

$$c_i(0) = \vec{m}_0 \cdot \vec{u}_i$$

$$f_i(t) = \vec{f}(t) \cdot \vec{u}_i$$



$$\frac{d\vec{u}}{dt} = \bar{A} \cdot \vec{u} + \int \vec{f} \quad \text{with} \quad \vec{u}(0) = \vec{u}_0$$

DRIVING  
FORCE  
(SOURCE)

I.C.

• Solution by LAPLACE and eigen decomposition of  $\bar{A}$ :

$$\vec{u}(t) \xrightarrow{\text{Laplace}} \vec{U}(s)$$

$$s \vec{U}(s) - \vec{u}_0 = \bar{A} \cdot \vec{U}(s) + \vec{F}(s)$$

$$s \vec{U}(s) - \bar{A} \cdot \vec{U}(s) = \vec{u}_0 + \vec{F}(s)$$

$$\begin{array}{ccc} \downarrow & & \downarrow \\ \text{I.C.} & & \text{DRIVING SOURCE} \end{array}$$

(NOTE: I.C. are equivalent to a driving source as a Dirac delta pulse at  $t=0$ !)

Eigen decomposition of  $\bar{A}$ :

$$\bar{A} = \bar{U} \cdot \bar{\Lambda} \cdot \bar{U}^{-1}$$

where:  $\bar{U}$  is square with columns  $\vec{u}_i$  eigenvectors

$\bar{\Lambda}$  is diagonal with elements  $\lambda_i$  eigenvalues

$\bar{U}^{-1}$  is the inverse of  $\bar{U}$  ( $= \bar{U}^T$  when  $\bar{A}$  is symmetric)

Corresponding EIGENMODE decomposition of  $\vec{U}(s)$ :

$$\vec{U}(s) = \bar{U} \cdot \vec{D}(s)$$

Also:

$$\vec{u}_0 = \bar{U} \cdot \vec{D}_0 \quad \text{decomposition of I.C.} \quad (\vec{D}_0 = \bar{U}^{-1} \vec{u}_0)$$

$$\vec{F}(s) = \bar{U} \cdot \vec{\Psi}(s) \quad \text{decomposition of SOURCE} \quad (\vec{\Psi}(s) = \bar{U}^{-1} \cdot \vec{F}(s))$$

$$\Rightarrow s \bar{U} \cdot \vec{D}(s) - \bar{U} \bar{\Lambda} \bar{U}^{-1} \bar{U} \vec{D}(s) = \bar{U} \cdot \vec{D}_0 + \bar{U} \cdot \vec{\Psi}(s)$$

$$(s - \bar{\Lambda}) \cdot \vec{D}(s) = \vec{D}_0 + \vec{\Psi}(s)$$

Individual components:

$$(s - \lambda_i) D_i(s) = D_{0i} + \Psi_i(s)$$

$$D_i(s) = \frac{D_{0i}}{s - \lambda_i} + \frac{1}{s - \lambda_i} \cdot \Psi_i(s)$$

$$\begin{aligned} \Rightarrow \vec{U}(s) &= \bar{U} \cdot \vec{D}(s) = \sum_{i=1}^N D_i(s) \cdot \vec{U}_i \\ &= \sum_{i=1}^N \frac{D_{0i}}{s - \lambda_i} \cdot \vec{U}_i + \sum_{i=1}^N \frac{1}{s - \lambda_i} \Psi_i(s) \vec{U}_i \end{aligned}$$

Inverse Laplace:

$$\vec{u}(t) = \sum_{i=1}^N b_{oi} e^{\lambda_i t} \vec{u}_i + \sum_{i=1}^N \int_0^t f_i(\theta) e^{\lambda_i(t-\theta)} d\theta \vec{u}_i$$

where  $\psi_i(s) \xrightarrow{\text{inverse Laplace}} f_i(t)$ , or

$$\vec{f}(t) = \sum_{i=1}^N f_i(t) \vec{u}_i \quad \text{DRIVING SOURCE DECOMPOSITION}$$

Also, I.C.  $\vec{u}(0) = \vec{u}_0$ , or

$$\vec{u}_0 = \sum_{i=1}^N \underbrace{b_{oi}}_{=c_i(0)} \vec{u}_i \quad \text{I.C. DECOMPOSITION}$$

Same as solution by variation of coefficients!